

Fourier Series Approximation for Max Operation in Non-Gaussian and Quadratic Statistical Static Timing Analysis

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Abstract—The most challenging problem in the current block-based statistical static timing analysis (SSTA) is how to handle the max operation efficiently and accurately. Existing SSTA techniques suffer from limited modeling capability by using a linear delay model with Gaussian distribution, or have scalability problems due to expensive operations involved to handle non-Gaussian variation sources or nonlinear delays. To overcome these limitations, we propose efficient algorithms to handle the max operation in SSTA with both quadratic delay dependency and non-Gaussian variation sources simultaneously. Based on such algorithms, we develop an SSTA flow with quadratic delay model and non-Gaussian variation sources. All the atomic operations, max and add, are calculated efficiently via either closed-form formulas or low dimension (at most 2-D) lookup tables. We prove that the complexity of our algorithm is linear in both variation sources and circuit sizes, hence our algorithm scales well for large designs. Compared to Monte Carlo simulation for non-Gaussian variation sources and nonlinear delay models, our approach predicts the mean, standard deviation and 95% percentile point with less than 2% error, and the skewness with less than 10% error.

Index Terms—Process variation, statistical static timing analysis (SSTA), timing.

I. INTRODUCTION

AS continuous CMOS technology scaling, process variation has become a potential show-stopper if not appropriately handled [1]–[6]. Statistical static timing analysis (SSTA) has thus become the frontier research topic in recent years in combating such variation effects. There are two types of SSTA methods, path-based SSTA [7]–[10] and block-based SSTA [11]–[19]. The goal of SSTA is to parameterize timing characteristics of the timing graph as a function of the underlying sources of process parameters that are modeled as

random variables. By performing SSTA, designers can obtain the timing distribution (yield) and its sensitivity to various process parameters. Such information is of tremendous value for both timing sign-off and design optimization for robustness and high profit margins.

The block-based SSTA is the most efficient SSTA method in recent years. In block-based SSTA, there are two major atomic operations, max and add. The add operation is simple, however, the max operation is much more complex. How to perform the max operation is the hardest problem in block-based SSTA. Although many studies have been done on this in recent years, the problem is far from being solved completely. For example, [11] and [12] assumed that all variation sources are Gaussian and independent of each other. Based on a linear delay model, [12] proposed a linear-time algorithm for SSTA, in which both atomic operations, max and add, can be performed efficiently via the concept of *tightness probability*. Because all variation sources are assumed to be Gaussian, so is the delay distribution under the linear delay model.

Such a Gaussian assumption is, however, no longer tolerable as more complicated or large-scale variation sources are taken into account in the nanometer manufacturing regime. For example, via resistance is known to be non-Gaussian with asymmetric distribution [17]; and dopant concentration is more suitably modeled as a Poisson distribution [16]. In addition, the linear dependency of delay on the variation sources is also not accurate, especially when variation becomes large [20]. For example, gate delay is inherently a nonlinear function of channel length and V_{th} [13], [17], which are two common sources of variation. Similarly, interconnect delay is also a nonlinear function of interconnect geometries [13], [14], which vary because of chemical-mechanical polishing. These combined non-Gaussian and nonlinear variation effects invalidate the linear delay model with Gaussian assumption in the existing SSTA.

Recently, non-Gaussian variation sources were addressed in [16], where independent component analysis (ICA) was used to find a set of independent components (not necessary Gaussian) to approximate the correlated non-Gaussian random variables. To do this, however, a complicated moment matching algorithm has to be used to make those atomic statistical operations feasible. Reference [21] introduced a conditional linear approximation of the max operation and captured all types of correlation. However, both of the above works are still based on a linear delay model, which cannot capture the nonlinear dependency of

Manuscript received November 09, 2010; revised April 19, 2011; accepted May 05, 2011. Date of publication June 30, 2011; date of current version June 14, 2012.

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Digital Object Identifier 10.1109/TVLSI.2011.2157843

delays on process parameters. To capture these nonlinear dependency effects, [13] and [14] proposed to use a quadratic delay model for SSTA. But to contain the complexity, they had to assume that all variation sources must follow a Gaussian distribution, even though the delay D itself may not be Gaussian. To compute $\max(D_1, D_2)$, [13] first developed closed form formulas to compute the mean and variance of the quadratic form. It then treats D_1 and D_2 as a Gaussian distribution to obtain the tightness probability. There is, however, no justification on why the tightness probability formula developed for Gaussian distributions can be applied for non-Gaussian distributions. [14] tried to re-construct $\max(D_1, D_2)$ through moment matching. To obtain those moments, however, many expensive numerical integration (2-D) operations have to be applied.

There are some existing studies [15], [17]–[19] trying to handle both nonlinear and non-Gaussian effects simultaneously. However, [15] computes $\max(D_1, D_2)$ by regression based on Monte Carlo simulation, which obviously is slow; [17] computed the max operation through tightness probability while [19] applied moment matching to reconstruct the max. However, to do so, both had to resort to expensive multi-dimension numerical integration techniques. [17] requires multi-dimension numerical integration to calculate the tightness probability and the number of integration dimensions depends on the number of non-Gaussian variation sources. [19] applies two-dimensional numerical integration to obtain the joint moments between variation sources and max. It is more efficient than [15] and [17]. But 2-D numerical integration is still an expensive operation for block-based SSTA. Hence scalability of such methods to handle a large number of non-Gaussian variation sources is limited. [18] handled the atomic operations by approximating the gate delay using a set of orthogonal polynomials, which needs to be constructed for different variation distributions.

In this work, we propose a novel method to handle the max operation for quadratic delay model with non-Gaussian variation sources. Based on such model, we develop a nonlinear and non-Gaussian SSTA technique (n^2 SSTA). The major advantages of this technique are multi-fold: 1) both nonlinear dependency and non-Gaussian variation sources are handled simultaneously for timing analysis; 2) all statistical atomic operations, max and add, are performed efficiently via either closed-form formulas or low dimension (at most 2-D) lookup tables; 3) the complexity of the n^2 SSTA algorithm is linear in both number of variation sources and circuit sizes. Compared to Monte Carlo simulation for non-Gaussian variation sources and nonlinear delay models, our approach predicts the mean, standard deviation and 95% percentile point with less than 2% error, and the skewness with less than 10% error.

The rest of the paper is organized as follows. Section II presents our nonlinear and non-Gaussian delay modeling and the basic block-based SSTA flow. Section III discuss our algorithm of the max operation with quadratic delay model and non-Gaussian variation sources. Section IV simplifies the method to handle linear delay model. We present experiments in Section V, and concludes this paper in Section VI.

II. PRELIMINARIES AND MODELING

A. Quadratic Delay Modeling

In general, device or interconnect delays of a design are a complicated nonlinear function of the underlying process parameters and it can be described as

$$D = F(X_1, X_2, \dots, X_i, \dots) \quad (1)$$

where the process parameters (such as channel length and threshold voltage) are modeled as a random variable X_i . In reality, the exact form of function F is not known, and X_i are not necessarily Gaussian.

The simplest approximation is the first- and second-order Taylor expansion as shown as follows:

$$D \approx d_0 + \sum a_i X_i \quad (2)$$

$$D \approx d_0 + \sum a_i X_i + \sum b_i X_i^2 + \sum_{i \neq k} b_{i,k} X_i X_k \quad (3)$$

where d_0 is the nominal value of D ; a_i and b_i are the first- and second-order sensitivities of D to X_i , respectively; and $b_{i,k}$ are the sensitivity to the joint variation of X_i and X_k . (2) is called the *first-order canonical form*, and is widely used for SSTA [11], [12]; whereas (3) is called the *quadratic delay model*, and has been studied in [13]–[15], [20]. But Gaussian assumptions limit their modeling capability and prevent them from reflecting the reality.

Therefore, we propose a different quadratic model as follows:

$$D = d_0 + \sum (a_i X_i + b_i X_i^2) + rR \quad (4)$$

where X_i represents global sources of variation, and R represents purely independent random variation which is modeled as a Gaussian random variable. Unlike previous work, we allow X_i to follow arbitrary random distributions. We refer to the delay model (4) as *general canonical form* in this paper.

B. Block-Based SSTA Framework

To compute the arrival time and required arrival time in a block-based SSTA framework, four atomic operations are sufficient, i.e., addition, subtraction, maximum, and minimum. Because of the symmetry between addition and subtraction (similarly maximum and minimum) operations, in the following, we will only discuss operations on addition and maximum. Given D_1 and D_2 in the form of (4)

$$D_1 = d_{10} + \sum (a_{1i} X_i + b_{1i} X_i^2) + r_1 R_1 \quad (5)$$

$$D_2 = d_{20} + \sum (a_{2i} X_i + b_{2i} X_i^2) + r_2 R_2 \quad (6)$$

we want to compute the $D_s = D_1 + D_2$ and $D_m = \max(D_1, D_2)$ in the following form:

$$D_s = d_{s0} + \sum (a_{si} X_i + b_{si} X_i^2) + r_s R_s \quad (7)$$

$$D_m = d_{m0} + \sum (a_{mi} X_i + b_{mi} X_i^2) + r_m R_m. \quad (8)$$

The add operation is straight forward. The coefficients of D_s can be computed by adding the correspondent coefficients of D_1 and D_2

$$d_{s0} = d_{01} + d_{02}; \quad a_{si} = a_{1i} + a_{2i} \quad (9)$$

$$b_{si} = b_{1i} + b_{2i}; \quad r_s = \sqrt{r_1^2 + r_2^2}. \quad (10)$$

The max operation is much complexer than the add operation, and we will discussed this in detail in Section III.

III. MAX FOR QUADRATIC DELAY MODEL

The max operation is the hardest operation for block-based SSTA. In this section, we propose a novel technique to efficiently compute the max of two general canonical forms, i.e., $D_m = \max(D_1, D_2)$. Without loss of the generality, in the following of this paper, we assume that $E[D_1] > E[D_2]$. We first compute the mean and variance of $V = D_1 - D_2$, if $\mu_V > 3\sigma_V$, then $D_m = D_1$. Otherwise, we compute the skewness of V and $E[\max(V, 0)]$. After that, we approximate the joint PDF of V and X_i 's with Fourier Series and compute the joint moments between V and X_i 's. Finally, we compute the joint moments between D_m and X_i 's and reconstruct the canonical form of D_m . In the following, we discuss our approach in details.

A. Moments of V

From the definition of V as discussed in the previous section, the max operation can be rewritten as

$$\max(D_1, D_2) = \max(D_1 - D_2, 0) + D_2 = \max(V, 0) + D_2. \quad (11)$$

According to (5) and (6), V can be written as the following form:

$$V = d_0 + \sum (a_i X_i + b_i X_i^2) + r_1 R_1 - r_2 R_2 \quad (12)$$

where $a_i = a_{1i} - a_{2i}$, $b_i = b_{1i} - b_{2i}$, and $d_0 = d_{01} - d_{02}$. In order to compute the central moments of V , we first rewrite V to the following form:

$$V = d'_0 + \sum_{i=1}^n Z_i + r_1 R_1 - r_2 R_2 \quad (13)$$

where

$$d'_0 = d_0 + \sum_{i=1}^n b_i \quad (14)$$

$$Z_i = a_i X_i + b_i X_i^2 - b_i. \quad (15)$$

From the above equations, it is easy to compute the first three moments of Z_i 's

$$E[Z_i] = 0 \quad (16)$$

$$E[Z_i^2] = b_i^2 (m_i(4) - 1) + 2a_i b_i m_i(3) + a_i^2 \quad (17)$$

$$E[Z_i^3] = 3a_i b_i^2 (m_i(5) - 2m_i(3)) + 3a_i^2 b_i (m_i(4) - 1) + a_i^3 m_i(3) + b_i^3 (m_i(6) - 3m_i(4) + 2). \quad (18)$$

Because X_i 's are mutually independent random variables, Z_i 's are independent random variables with zero mean. Therefore, the first three central moments of V are

$$\mu_V = E[V] = d'_0 \quad (19)$$

$$\sigma_V^2 = E[(V - \mu_V)^2] = r_1^2 + r_2^2 + \sum_{i=1}^n E[Z_i^2] \quad (20)$$

$$m_V(3) = E[(V - \mu_V)^3] = r_1^3 + r_2^3 + \sum_{i=1}^n E[Z_i^3]. \quad (21)$$

With the central moments, the skewness of V can be computed as

$$\gamma_V = m_V(3)/\sigma_V^3. \quad (22)$$

After computing the mean and variation of V , we can quickly compute some trivial max operation. When μ_V is larger than the $m\text{-}\sigma$ value, that is, $\mu_V > m \cdot \sigma_V$, which means the probability that $D_1 > D_2$ is very high, we may let $\max(D_1, D_2) \approx D_1$. The range value m can be set by the users, in the following of this paper, we let $m = 3$.

B. Mean of the Max Operation

In order to reconstruct the canonical form of D_m , we first need to compute $\mu_{V_m} = E[\max(V, 0)]$. When V is a non-Gaussian random variable, exact computation of μ_{V_m} is difficult in general. Therefore, we propose to use the following two-step procedure to approximately compute μ_{V_m} . In the first step, we approximate the non-Gaussian random variable V as a quadratic function of a standard Gaussian random variable W similar to [18] and [24], i.e.,

$$V \approx P(W) = c_2 \cdot W^2 + c_1 \cdot W + c_0. \quad (23)$$

The coefficients c_2 , c_1 , and c_0 can be obtained by matching $P(W)$ and V 's mean, variance, and skewness simultaneously, i.e.,

$$\mu_V = E[c_2 \cdot W^2 + c_1 \cdot W + c_0] = c_2 + c_0 \quad (24)$$

$$\sigma_V^2 = E[(c_2 \cdot W^2 + c_1 \cdot W + c_0 - \mu_V)^2] = 2c_2^2 + c_1^2 \quad (25)$$

$$\gamma_V \cdot \sigma_V^3 = E[(c_2 \cdot W^2 + c_1 \cdot W + c_0 - \mu_V)^3] = 8c_2^3 + 6c_2 c_1^2. \quad (26)$$

As shown in [24], solving the above equations, c_2 will be one of the following values:

$$c_{2,1} = -\frac{2\sigma_V^2 + \Delta^{2/3}}{2\Delta^{1/3}} \quad (27)$$

$$c_{2,2} = \frac{2(1 + j\sqrt{3})\sigma_V^2 + (1 - j\sqrt{3})\Delta^{2/3}}{4\Delta^{1/3}} \quad (28)$$

$$c_{2,3} = \frac{2(1 - j\sqrt{3})\sigma_V^2 + (1 + j\sqrt{3})\Delta^{2/3}}{4\Delta^{1/3}} \quad (29)$$

where $\Delta = \gamma_V \cdot \sigma_V^3 + j\sqrt{8\sigma_V^6 - \gamma_V^2 \cdot \sigma_V^6}$, with $j = \sqrt{-1}$. The mean, variance, and skewness of V can be computed from (19), (20), and (22), respectively. It is proved in [24] that, when $|\gamma_V| < 2\sqrt{2}$ there exists one and only one of the above three

values in the range $|c_2| < \sigma_V/\sqrt{2}$. We will pick such value for c_2 . When $|\gamma_V| > 2\sqrt{2}$, we may let

$$c_2 = \begin{cases} \sigma_V/\sqrt{2} & \gamma_V > 2\sqrt{2} \\ -\sigma_V/\sqrt{2} & \gamma_V < -2\sqrt{2}. \end{cases} \quad (30)$$

It is proved in [18] that the above equations gives $P(W)$ which matches the mean and variance of V and has the skewness closest to γ_V . With c_2 , c_1 , and c_0 can be computed as

$$c_1 = \sqrt{\sigma_V^2 - 2c_2^2} \quad (31)$$

$$c_0 = \mu_V - c_2. \quad (32)$$

After obtaining the coefficients c_2 , c_1 , and c_0 in the second step, we approximate the exact mean of $\max(V, 0)$ by the exact mean of $\max(P(W), 0)$, i.e.,

$$\mu_{Vm} \approx E[\max(P(W), 0)] = \int_{P(w)>0} P(w)\phi(w)dw \quad (33)$$

where $\phi(\cdot)$ is the PDF of the standard normal distribution. In the above equation, the integration range $P(w) > 0$ can be computed in the following four different cases.

Case 1) $c_2 > 0$

$$P(w) > 0 \Leftrightarrow w < t_1 \vee w > t_2 \quad (34)$$

where

$$t_1 = \left(-c_1 - \sqrt{c_1^2 - 4c_2c_0}\right)/2c_2 \quad (35)$$

$$t_2 = \left(-c_1 + \sqrt{c_1^2 - 4c_2c_0}\right)/2c_2. \quad (36)$$

Case 2) $c_2 < 0$

$$P(w) > 0 \Leftrightarrow t_2 < w < t_1. \quad (37)$$

Case 3) $c_2 = 0 \wedge c_1 > 0$

$$P(w) > 0 \Leftrightarrow w > -c_0/c_1. \quad (38)$$

Case 4) $c_2 = 0 \wedge c_1 < 0$

$$P(w) > 0 \Leftrightarrow w < -c_0/c_1. \quad (39)$$

Knowing the integration range, we can compute $E[\max(P(W), 0)]$ under such four cases

$$E[\max(P(W), 0)] = \begin{cases} (c_2 + c_0)(1 + \Phi(t_1) - \Phi(t_2)) \\ \quad + (c_1 + t_1)(\phi(t_2) - \phi(t_1)) & c_2 > 0 \\ (c_2 + c_0)(\Phi(t_1) - \Phi(t_2)) \\ \quad + (c_1 + t_1)(\phi(t_2) - \phi(t_1)) & c_2 < 0 \\ c_0 \cdot \Phi(c_0/c_1) + c_1 \cdot \phi(c_0/c_1) & c_2 = 0 \wedge c_1 > 0 \\ c_0 \cdot (1 - \Phi(c_0/c_1)) - c_1 \cdot \phi(c_0/c_1) & c_2 = 0 \wedge c_1 < 0 \end{cases} \quad (40)$$

where $\Phi(\cdot)$ is the cumulative density function (CDF) of the standard normal distribution. According to (41), we can compute μ_{Vm} easily through analytical formulas.

C. Joint Moments Between $\max(V, 0)$ and Variation Sources

In the previous section, we compute the mean of $\max(V, 0)$. In order to reconstruct the canonical form of D_m , we also need to know the joint moments between V and the variation sources X_i 's. Because V and X_i 's are correlated non-Gaussian random variables, the computation of the joint moments between them is complex. In this section, we will introduce a new method based on Fourier Series expansion to solve such problem.

In order to compute the joint moments, we first approximate the joint PDF of V and the i th variation sources X_i , $f_i(v, x_i)$, by its first K th order 2-D Fourier Series within the 3σ range

$$f_i(v, x_i) \approx \sum_{p,q=-K}^K \alpha_{i,pq} \cdot e^{\zeta_p v + \eta_{i,q} x_i} \quad (41)$$

where $\zeta_p = jp\pi/l$ and $\eta_{i,q} = jq\pi/h_i$, with $l = 3\sigma_V$, and $h_i = 3\sigma_{X_i}$. $\alpha_{i,pq}$ are the Fourier coefficients and they can be computed as

$$\alpha_{i,pq} = \frac{1}{4lh_i} \int_{\mu_V-l-h}^{\mu_V+l-h} \int e^{-\zeta_p v_1 - \eta_{i,q} v_2} \cdot f_i(v, x_i) dv dx_i. \quad (42)$$

Considering that both V and X_i are within the 3σ range, i.e., $f_i(v, x_i) \approx 0$ when $v \notin (\mu_V - l, \mu_V + l)$ or $x_i \notin (\mu_{X_i} - h_i, \mu_{X_i} + h_i)$, (43) can be further simplified as

$$\begin{aligned} \alpha_{i,pq} &= \frac{1}{4lh_i} E[e^{-\zeta_p V - \eta_{i,q} X_i}] \\ &= \frac{1}{4lh_i} e^{-Y_{c,p}} E\left[e^{-Y_{r1,p} - Y_{r2,p} - \widehat{Y}_{i,pq} - \sum_{j \neq i} Y_{j,p}}\right] \end{aligned} \quad (43)$$

where

$$Y_{c,p} = \zeta_p \cdot d_0 \quad (44)$$

$$Y_{r1,p} = \zeta_p \cdot r_1 R_1 \quad (45)$$

$$Y_{r2,p} = \zeta_p \cdot r_2 R_2 \quad (46)$$

$$Y_{i,p} = \zeta_p (a_i X_i + b_i X_i^2) \quad (47)$$

$$\widehat{Y}_{i,pq} = (\zeta_p a_i + \eta_{i,q}) X_i + \zeta_p b_i X_i^2. \quad (48)$$

Because all X_i 's are independent, so are all $Y_{j,p}$'s, $Y_{r1,p}$, $Y_{r2,p}$, and $\widehat{Y}_{i,pq}$. Then $\alpha_{i,pq}$ can be further simplified as

$$\begin{aligned} \alpha_{i,pq} &= \frac{1}{4lh_i} e^{-Y_{c,p}} \cdot E[e^{-Y_{r1,p}}] \cdot E[e^{-Y_{r2,p}}] \\ &\quad \cdot E[e^{-\widehat{Y}_{i,pq}}] \cdot \prod_{k \neq i} E[e^{-Y_{k,p}}]. \end{aligned} \quad (49)$$

As $\widehat{Y}_{i,pq}$, $Y_{r1,p}$, $Y_{r2,p}$, and $Y_{k,p}$'s can be written as a general form as

$$Y = s_1 X_i + s_2 X_i^2 \quad (50)$$

TABLE I
EXPERIMENT SETTING TO VERIFY $\max(V, X_i)$

	d_0	a_i	b_i	r
V	0	{4,3,4,4}	{1,2,1,1}	1

with s_1 and s_2 being two constant values, in the following, we discuss how to compute $E[e^{-Y}]$ in its general form. By definition

$$E[e^{-Y}] = \int_{-h_i}^{h_i} e^{-s_1 x_i - s_2 x_i^2} g_i(x_i) dx_i \quad (51)$$

where $g_i(x_i)$ is PDF of X_i , whose range is given by $-h_i \leq X_i \leq h_i$. For arbitrary $g_i(x_i)$, we can build a 2-D table indexed by s_1 and s_2 to compute (51).

To validate our computation of JPDF of V and X_i , we compare our computed JPDF with Monte Carlo simulated JPDF. One of the examples is shown in Table I with four sources of random variables (i.e., X_i for $i = 1, 2, 3, 4$) that all follow a uniform distribution. In the experiment, we compute the JPDF of V and X_1 . The order of Fourier series to approximate JPDF is four ($K = 4$).

With the Fourier Series approximation discussed above, we can compute the first two joint moments between V and X_i . By definition, the joint moments can be computed as

$$\begin{aligned} Ec_{i,1} &= E[X_i \cdot \max(V, 0)] \\ &= \int_{-h_i}^{h_i} \int_0^{l+\mu_V} vx_i \cdot f_i(v, x_i) dv dx_i \end{aligned} \quad (52)$$

$$\begin{aligned} Ec_{i,2} &= E[X_i^2 \cdot \max(V, 0)] \\ &= \int_{-h_i}^{h_i} \int_0^{l+\mu_V} vx_i^2 \cdot f_i(v, x_i) dv dx_i. \end{aligned} \quad (53)$$

Replacing the joint PDF $f_i(v, x_i)$ with its Fourier Series approximation, we have

$$\begin{aligned} Ec_{i,1} &= \int_{-h_i}^{h_i} \int_0^{l+\mu_V} vx_i \cdot \sum_{p,q=-K}^K \alpha_{i,pq} \cdot e^{\zeta_p v + \eta_{i,q} x_i} dv dx_i \\ &= \sum_{p,q=-K}^K \alpha_{i,pq} \cdot U_p \cdot T_{iq,1} \end{aligned} \quad (54)$$

$$\begin{aligned} Ec_{i,2} &= \int_{-h_i}^{h_i} \int_0^{l+\mu_V} vx_i^2 \cdot \sum_{p,q=-K}^K \alpha_{i,pq} \cdot e^{\zeta_p v + \eta_{i,q} x_i} dv dx_i \\ &= \sum_{p,q=-K}^K \alpha_{i,pq} \cdot U_p \cdot T_{iq,2} \end{aligned} \quad (55)$$

where U_p , $T_{iq,1}$, and $T_{iq,2}$ can be computed as

$$\begin{aligned} U_p &= \int_0^{l+\mu_V} v e^{\zeta_p v} dv \\ &= \begin{cases} \frac{1}{\eta_p^2} ((-1)^p e^{\zeta_p \mu_V} (\zeta_p \mu_V + \zeta_p l - 1) + 1) & p \neq 0 \\ \frac{(l+\mu_V)^2}{2} & p = 0 \end{cases} \end{aligned} \quad (56)$$

$$T_{iq,1} = \int_{-h_i}^{h_i} x_i e^{\eta_{i,q} x_i} dx_i = \begin{cases} (-1)^q \frac{2h_i^2}{jn\pi} & q \neq 0 \\ 0 & q = 0 \end{cases} \quad (57)$$

$$T_{iq,2} = \int_{-h_i}^{h_i} x_i^2 e^{\eta_{i,q} x_i} dx_i = \begin{cases} (-1)^q \frac{4h_i^3}{n^2\pi^2} & q \neq 0 \\ \frac{2}{3} h_i^3 & q = 0. \end{cases} \quad (58)$$

From the above discussion, we see that the joint moments between V and X_i 's can be computed by close form formulas. Moreover, the joint moments between $\max(V, 0)$ and random variation sources, $E[R_1 \max(V, 0)]$, $E[R_2 \max(V, 0)]$ can be computed in the same way.

D. Reconstruction of the Canonical Form

With mean of $\max(V, 0)$ and the joint moments between $\max(V, 0)$ and X_i 's computed in the previous section, it is easy to compute the mean of D_m and the joint moments between D_m and X_i 's

$$\begin{aligned} E[D_m] &= E[\max(V, 0)] + E[D_2] \\ &= \mu_{Vm} + d_{02} + \sum b_{2i} \end{aligned} \quad (59)$$

$$\begin{aligned} E[X_i D_m] &= E[X_i (\max(V, 0) + D_2)] \\ &= E[X_i \max(V, 0)] + E[X_i D_2] \\ &= Ec_{i,1} + a_{2i} + b_{2i} \cdot m_i(3) \end{aligned} \quad (60)$$

$$\begin{aligned} E[X_i^2 D_m] &= E[X_i^2 \max(V, 0)] + E[X_i^2 D_2] \\ &= Ec_{i,2} + a_{2i} \cdot m_i(3) + b_{2i} \cdot m_i(4) \\ &\quad + \sum_{k \neq i} b_{2k} + d_{02} \end{aligned} \quad (61)$$

$$\begin{aligned} E[R_1 D_m] &= E[X_i \max(V, 0)] + E[R_1 D_2] \\ &= E[R_1 \max(V, 0)] \end{aligned} \quad (62)$$

$$\begin{aligned} E[R_2 D_m] &= E[X_i \max(V, 0)] + E[R_2 D_2] \\ &= E[R_2 \max(V, 0)] + r_2. \end{aligned} \quad (63)$$

Because we want to reconstruct D_m in the second-order canonical form, as shown in (8), by applying the moment matching technique similar to [14], we have

$$E[D_m] = \sum_{i=1}^n b_{mi} + d_{m0} \quad (64)$$

$$E[X_i D_m] = a_{mi} + b_{mi} \cdot m_i(3) \quad (65)$$

$$\begin{aligned} E[X_i^2 D_m] &= a_{mi} \cdot m_i(3) + b_{mi} \cdot m_i(4) \\ &\quad + \left(\sum_{k \neq i} b_{mk} + d_{m0} \right) \\ &= a_{mi} \cdot m_i(3) + b_{mi} \cdot (m_i(4) - 1) \\ &\quad + E[D_m]. \end{aligned} \quad (66)$$

With the joint moments computed in (59)–(61), the constant term d_{m0} and the sensitivity coefficients a_{mi} 's and b_{mi} 's can be obtained by solving the linear equations above

$$b_{mi} = \frac{E[X_i^2 D_m] - E[D_m] - E[X_i D_m] m_i(3)}{m_i(4) - m_i^2(3) - 1} \quad (67)$$

$$a_{mi} = E[X_i D_m] - b_{mi} \cdot m_i(3) \quad (68)$$

$$d_{m0} = E[D_m] - \sum_{i=1}^n b_{mi}. \quad (69)$$

Because the random term in D_m comes from the random terms in D_1 and D_2 , we assume that $r_m R_m = r_{m1} R_1 + r_{m2} R_2$. Because the random variation sources R_m , R_1 , and R_2 are Gaussian random variables, by applying the moment matching technique similar to (67), we have

$$r_{m1} = E[R_1 \cdot D_m] \quad (70)$$

$$r_{m2} = E[R_2 \cdot D_m] \quad (71)$$

$$r_m = \sqrt{r_{m1}^2 + r_{m2}^2} \quad (72)$$

where $E[R_1 \cdot D_m]$ and $E[R_2 \cdot D_m]$ are computed in (63) and (64), respectively.

E. Computational Complexity Analysis

For each max operation, we need to compute n joint moments between X_i 's and D_m , where n is the total number of variation sources. For each variation source, we need to compute $4(K+1)^2$ Fourier coefficients for the joint PDF approximation, where K is the maximum order of Fourier Series. Therefore, the total computational complexity of one step max operation is $\mathcal{O}\{nK^2\}$. Usually, larger K provides more accurate approximation. Our experiment shows that $K = 4$ provides very accurate approximation for the joint PDF. For the add operation, it is easy to see that the computational complexity is $\mathcal{O}\{n\}$. The total number of max and add operations for block-based SSTA is linear to the circuit size.

Moreover, in our original algorithm for n^2 SSTA [25], the computational complexity for one step max operation is $\mathcal{O}\{nK^3\}$. In this paper, we have improved the algorithm and reduced the complexity to $\mathcal{O}\{nK^2\}$. The experimental results in Section V show that the improved n^2 SSTA is much faster than our original algorithm in [25]. In addition, because that our method assumes arbitrary distribution for variation sources, the distribution of variation sources does not affect the efficiency of our method, no matter the variation sources are with Gaussian distribution or not.

IV. MAX OPERATION FOR LINEAR DELAY MODEL

In the previous sections, we introduce the SSTA for nonlinear delay model. In practice, when the variation scale is small, the circuit delay can be approximated as a linear function of variation sources

$$D = d_0 + \sum a_i X_i + r \cdot R \quad (73)$$

where X , d_0 , A , r , and R are defined in the similar way as the second-order delay model in (4). The difference is that there are no second-order terms. Hence the atomic operations of the linear delay model are simpler than those of quadratic delay model.

For the linear delay model, we can compute the mean of D_m in the similar way as in (40) and (59). For the joint moments, we can also apply the Fourier Series approximation method as shown in Section III-C. The only difference is that when computing the Fourier coefficients, Y s in (50) is in the form of $Y = cX$. Hence $E[e^Y] = \mathcal{F}_i(-c/j)$, where $\mathcal{F}_i(\cdot)$ is the Fourier transformation of $g_i(x_i)$. $\mathcal{F}_i(\cdot)$ can be precomputed and stored in a 1-D lookup table. Such computation is simpler than the second order delay model. After computing the joint moments,

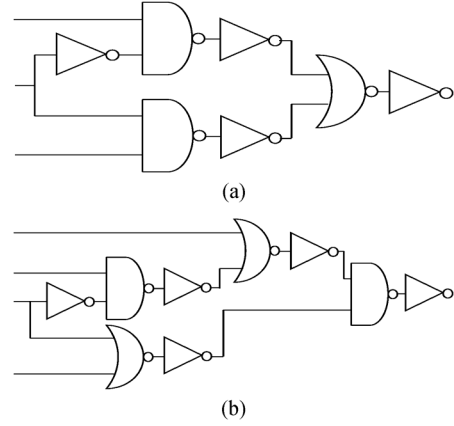


Fig. 1. Test circuits. (a) Circuit 1. (b) Circuit 2.

we can reconstruct the linear form of D_m in the same way as the second-order delay model as shown in Section III-D. Moreover, it is easy to find that the computational complexity of the SSTA for linear delay model is the same as that for second-order delay model.

V. EXPERIMENTAL RESULT

We have implemented our SSTA algorithm in C for both second-order delay model (n^2 SSTA) and linear delay model (*Lin-SSTA*). In our experiment, we assume that the maximum order of Fourier Series approximation $K = 4$ for both delay models. For comparison, we also define two comparison cases: 1) our implementation of the linear SSTA for Gaussian variation sources in [12], which we refer to as *Lin-Gau* and 2) 100 000-sample Monte-Carlo simulation (*MC*). We apply all the above methods to the ISCAS89 suite of benchmarks in TSMC 65-nm technology.

In our experiment, we consider two types of variation sources L_{eff} and V_{th} . For each type of variation sources, inter-die, intra-die spatial, and intra-die random variations are considered. We modeled the spatial variation using the grid-based model in [26]. The number of grids (the number of spatial variation sources) is determined by the circuit size, larger circuit will have more variation sources. We also assume that the 3σ value of the inter-die, intra-die spatial, and intra-die random variation are 10%, 10%, and 5% of the nominal value, respectively. In the following, We perform the experiments for two variation setting: 1) both L_{eff} and V_{th} are with skew-normal distribution [27] and 2) L_{eff} is with normal distribution and V_{th} is with Poisson distribution. The experimental setting is shown in Table II.

In Section II, we have verified the accuracy of quadratic fitting of gate delay. In this section, we further verify the accuracy of Monte Carlo simulation with quadratic delay fitting. Because it is very time consuming to perform Monte Carlo SPICE simulation on big benchmarks, we construct two small test circuits as shown in Fig. 1. Table V compares the mean, standard deviation, and skewness for the 10 000-sample SPICE Monte Carlo simulation, 10 000-Monte Carlo simulation with quadratic delay fitting, and n^2 SSTA under variation setting (1). From the table, we see that the result of Monte Carlo simulation with quadratic

TABLE II
EXPERIMENT SETTING. THE 3σ VALUE IS THE NORMALIZED WITH RESPECT TO THE NOMINAL VALUE

	L_{eff} distribution	V_{th} distribution	3σ inter-die	3σ spatial	3σ random
Case (1)	Skewnormal	Skewnormal	10%	10%	5%
Case (2)	Normal	Poisson	10%	10%	5%

TABLE III
COMPARISON BETWEEN SPICE SIMULATION, MONTE CARLO SIMULATION AND n^2 SSTA

bench mark	SPICE MC			MC with quad fitting			n^2 SSTA		
	μ (ps)	σ (ps)	γ	μ (ps)	σ (ps)	γ	μ (ps)	σ (ps)	γ
circuit1	541.2	98.3	0.52	543.4 (+0.41%)	98.9 (+0.61%)	0.54 (+3.8%)	544.2 (+0.55%)	99.1 (+0.8%)	0.55 (+5.8%)
circuit2	792.2	124.3	0.59	795.5 (+0.4%)	126.2 (+1.5%)	0.61 (+3.3%)	796.1 (+0.5%)	127.1 (+2.3%)	0.64 (+7.9%)

TABLE IV
ERROR PERCENTAGE OF MEAN, STANDARD DEVIATION, SKEWNESS, AND 95% PERCENTILE POINT FOR VARIATION SETTING (1). NOTE: THE ERROR PERCENTAGE OF MEAN, STANDARD DEVIATION, AND 95% PERCENTILE POINT IS COMPUTED AS $100 \times (MC_value - SSTA_value) / \sigma_{MC}$, AND THE ERROR PERCENTAGE OF SKEWNESS IS COMPUTED AS $100 \times (\gamma_{MC} - \gamma_{SSTA}) / \gamma_{MC}$

bench name	G	N	n^2 SSTA					Lin-SSTA					Lin-Gau					[25]	MC
			μ	σ	γ	95%	T	μ	σ	γ	95%	T	μ	σ	γ	95%	T		
s27	8	8	0.25	-0.71	7.34	1.23	0.65	0.41	3.45	-27.23	-4.67	0.32	-1.23	-4.06	-100	-23.14	0.01	1.12	1.71
s344	101	18	0.63	-1.18	-2.24	1.21	1.13	1.02	-1.66	-17.35	-3.89	0.61	-1.10	1.52	-100	-20.82	0.02	5.54	34.29
s386	118	18	0.67	-1.40	-3.00	1.04	1.18	0.53	-2.41	-15.94	6.50	0.65	1.07	-1.43	-100	-21.82	0.02	6.49	36.03
s420	140	25	0.69	-1.25	-3.21	1.43	1.36	-0.75	1.71	-18.79	-4.06	0.69	1.33	1.24	-100	-25.30	0.02	7.35	39.06
s444	119	25	-0.61	-1.06	-4.34	-1.16	1.38	-1.31	2.06	-24.07	-6.18	0.67	-1.33	1.28	-100	-29.26	0.02	8.69	41.10
s832	262	33	-0.31	-1.08	6.91	-1.47	1.92	0.28	-4.99	-27.40	-2.82	0.85	1.14	-4.71	-100	-16.47	0.04	10.25	102.7
s953	311	43	0.78	-1.15	-2.36	1.63	1.96	-0.84	2.09	-17.18	-5.65	0.87	1.29	1.18	-100	-25.36	0.05	11.32	155.7
s1196	388	55	-0.30	-0.69	8.17	-1.04	2.17	-0.30	-4.89	-35.69	2.20	0.96	1.31	-4.40	-100	-21.79	0.05	13.13	172.2
s1494	588	68	-0.78	-1.30	-3.69	-1.77	2.27	1.20	1.75	-16.15	5.88	1.03	-1.05	1.16	-100	-28.69	0.06	15.52	323.4
s5378	1004	82	-0.65	-1.13	-2.67	-1.65	5.37	0.73	2.22	-22.54	6.14	2.26	-1.00	1.49	-100	-30.09	0.07	27.32	1256
s9234	2027	99	-1.18	-0.18	2.85	-0.64	6.83	0.65	-1.69	-17.37	4.15	3.14	-3.07	-1.47	-100	-31.81	0.09	32.18	2769
s13207	2573	115	-0.93	-0.25	2.30	-0.52	9.02	0.60	1.75	-21.27	5.74	4.22	2.29	-1.86	-100	-26.34	0.11	47.35	4715
s15850	3448	135	-1.51	-0.21	2.87	-0.51	11.36	0.70	2.32	-20.45	3.44	5.28	3.35	1.33	-100	-33.16	0.15	58.25	6891
s38417	8709	176	-0.25	-0.91	6.14	-1.03	23.37	-2.14	3.53	-30.24	-3.38	10.00	-1.11	-5.30	-100	-16.58	0.28	98.13	19061
s38584	11448	176	-1.41	-0.27	2.42	-0.55	25.39	-1.95	-0.90	-19.85	3.96	13.49	-2.40	-1.82	-100	-28.40	0.31	104.1	19203
Ave	-	-	0.73	0.85	4.03	1.13	-	0.89	2.49	22.15	4.58	-	1.60	2.58	100	25.30	-	-	-

TABLE V
ERROR PERCENTAGE OF MEAN, STANDARD DEVIATION, SKEWNESS, AND 95% PERCENTILE POINT FOR VARIATION SETTING (2)

bench name	G	N	n^2 SSTA					Lin-SSTA					Lin-Gau					[25]	MC
			μ	σ	γ	95%	T	μ	σ	γ	95%	T	μ	σ	γ	95%	T		
s27	8	8	0.29	-0.93	8.21	1.07	0.64	0.37	-3.39	-30.01	-4.06	0.32	-1.24	-4.15	-100	-22.71	0.01	1.11	1.74
s344	101	18	-0.66	-1.40	-2.29	-1.26	1.12	-1.20	-1.91	-19.22	4.17	0.61	-1.12	1.55	-100	-21.58	0.02	5.34	33.84
s386	118	18	-0.65	-1.36	-2.93	-1.20	1.17	0.60	1.71	-19.13	6.65	0.64	1.08	-1.44	-100	-22.72	0.02	6.52	35.92
s420	140	25	-0.61	-1.28	-4.50	-1.45	1.38	-0.63	-1.71	-20.04	-5.44	0.66	1.30	1.19	-100	-24.26	0.02	7.33	38.45
s444	119	25	-0.87	-1.52	-3.20	-1.19	1.38	-1.26	2.12	-23.03	-7.41	0.69	-1.43	1.30	-100	-29.81	0.02	8.74	40.36
s832	262	33	0.27	-0.90	8.75	1.30	1.87	-0.26	4.86	-30.64	3.24	0.84	1.13	-4.63	-100	-16.88	0.04	10.39	99.4
s953	311	43	0.79	-1.16	-3.30	1.34	1.92	-0.68	-1.68	-14.91	-5.19	0.90	1.25	1.09	-100	-25.09	0.05	11.46	152.6
s1196	388	55	0.33	-0.69	6.32	1.38	2.13	-0.24	4.59	-31.14	2.75	0.93	1.39	-4.02	-100	-21.74	0.05	13.17	173.9
s1494	588	68	-0.85	-1.55	-3.46	-1.82	2.37	1.48	-2.54	-14.35	7.10	1.03	-1.03	1.21	-100	-26.44	0.06	15.78	317.6
s5378	1004	82	0.71	-1.20	-3.61	1.76	5.38	-0.74	2.38	-20.69	-5.59	2.34	-1.02	1.45	-100	-28.92	0.07	27.12	1232
s9234	2027	99	1.05	-0.23	2.22	0.53	6.77	0.61	-1.64	-16.17	4.52	3.15	-3.24	-1.39	-100	-31.05	0.09	32.60	2804
s13207	2573	115	1.16	-0.25	2.49	0.72	8.83	-0.56	-1.83	-18.04	-5.53	4.17	2.40	-1.84	-100	-26.56	0.11	48.05	4639
s15850	3448	135	1.22	-0.24	2.95	0.51	11.40	0.79	-1.72	-19.77	3.64	5.16	3.34	1.35	-100	-33.65	0.15	57.15	6653
s38417	8709	176	0.25	-0.70	6.17	1.20	23.23	1.76	3.54	-27.25	3.47	10.28	-1.15	-5.11	-100	-16.60	0.28	99.22	19470
s38584	11448	176	-1.56	-0.21	1.92	-0.54	25.28	2.17	0.97	-21.92	-4.13	13.39	-2.59	-1.72	-100	-29.39	0.31	105.0	19891
Ave	-	-	0.75	0.91	4.16	1.15	-	0.89	2.44	21.8	4.86	-	1.65	2.53	100	25.2	-	-	-

delay fitting is very close to the SPICE Monte Carlo simulation. In the rest of this section, we will use the Monte Carlo simulation with quadratic delay fitting as the golden case for comparison.

Fig. 2 illustrates the PDF comparison for circuit s15850 under variation setting (1). From the figure, we find that compared to the Monte-Carlo simulation, n^2 SSTA is the most accurate, the next is *Lin-SSTA*. Both of our SSTA methods are more accurate than *Lin-Gau*. Such result is reasonable because n^2 SSTA captures the both nonlinear and linear effects while *Lin-SSTA* captures only linear effect. *Lin-Gau* is for Gaussian variation

sources only, and therefore it does not work well under non-Gaussian variation sources.

Table V compares the run time in second (T), and the error percentage of mean (μ), standard deviation (σ), skewness (γ), and 95% percentile point (95%) under variation setting (1). In the table, the error percentage of mean, standard deviation and 95% percentile point is computed as $100 \times (MC_Value - SSTA_value) / \sigma_{MC}$, and the error percentage of skewness is computed as $100 \times (\gamma_{MC} - \gamma_{SSTA}) / \gamma_{MC}$. Moreover, the average error in the table is average of the absolute value. From

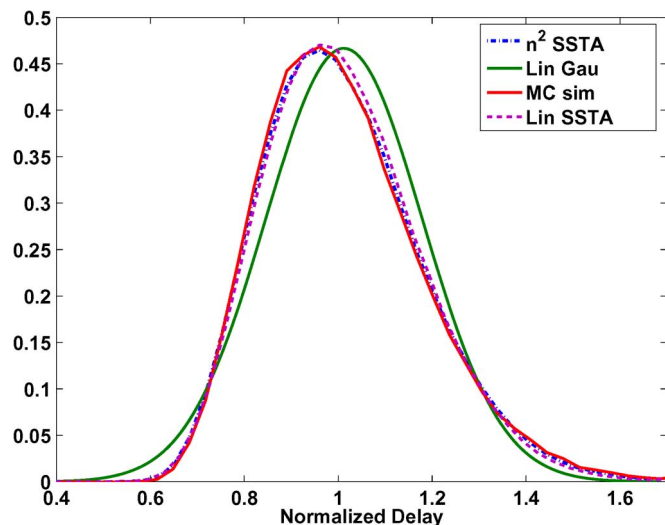


Fig. 2. PDF comparison for circuit s15850.

the table, we see that for n^2 SSTA the error of mean, standard deviation and 95% percentile point is within 2%, and the error of skewness is within 8%. The Lin -SSTA results similar mean deviation error. Compared to n^2 SSTA, the error of standard deviation and 95% percentile point is a little bit higher, but the error of skewness is much larger. This is because Lin -SSTA ignores all nonlinear effects which significantly affect the skewness. The error of Lin -Gau is larger than both of our SSTA methods, especially for skewness. This is because Lin -Gau is for Gaussian variation sources and cannot capture the skewness of the non-Gaussian variation sources. Moreover, we also find that the run time of all SSTA methods are linear to the circuit size. The run time of our methods is larger than that of Lin -Gau, but it is still acceptable and is significantly shorter than that of the Monte-Carlo simulation. In the table, we also compare the run time of our algorithm in this paper to that of our original method [25]. We see the approach in this paper is about $5\times$ faster than our original method. Because the basic algorithms of this paper and our original method in [25] are the same, the approach in this paper provides the same accuracy as our original method.

Table V illustrates the results under variation setting (2). It can be found that our methods are still very accurate under such variation setting. For the n^2 SSTA, the error of mean, standard deviation, and 95% percentile point is within 2%, and the error of skewness is within 10%. This shows that our approach works well for different distributions.

VI. CONCLUSION

A novel method to handle the max operation has been presented to handle both quadratic delay dependency and non-Gaussian variation sources simultaneously. An SSTA flow, n^2 SSTA, has been developed based on such max operation. We have shown that all statistical atomic operations can be performed efficiently via either closed-form formulas or low dimension (at most 2-D) lookup table. It has been proved that the complexity of n^2 SSTA is linear in both variation sources and circuit sizes. Compared to Monte Carlo simulation

for non-Gaussian variations and nonlinear delay models, our approach predicts the mean, standard deviation and 95% percentile point with less than 2% error, and the skewness with less than 10% error.

In the future, we will extend our work to consider more general delay models, such as non-polynomial delays and/or dependency on variations' cross terms. Moreover, in this paper, we use deterministic slew rate model. We will also try to consider statistical slew model in our future research.

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